# ON NUMERICAL REALIZATION OF THE FUNCTION OF THE HEREDITARY OPERATOR 

PMM Vol. 42, No. 6, 1978, pp. 1115-1122<br>E. S. SINAISKII<br>(Dnepropetrovsk)<br>(Received March 28, 1978)


#### Abstract

The problem of numerical realization of a function of the hereditary operator acting on some function of time is considered. Laplace transformations are used for the operators with kemels of the Rabotnov and Rzhanitsyn type to obtain formulas which reduce the problem in question to that of computing a quadrature. When the variable assumes large values, the formulas become asymptotic equations with an estimable error of approximation.


1. Effective solution of a wide class of problems of hereditary elasticity requires the application of the Volterra principle [1]. The replacement of the elastic constants of the material by the corresponding rheological operators, which is carried out in the solution of the problem for a perfectly elastic body, leads to the necessity of computing the convolutions of the form

$$
\begin{align*}
& \varphi\left(x ; q^{*}\right) f(t), \quad q^{*} f(t) \equiv \int_{0}^{t} q(t-\tau) f(\tau) d \tau  \tag{1.1}\\
& q_{j}^{*} q_{k}^{*} f=q_{j}^{*}\left(q_{k}^{*} f\right), \quad q^{*^{n}} f=q^{*}\left(q^{*^{n-1}} f\right)
\end{align*}
$$

Here $\varphi$ is a function of the spacial coordinates $x_{k}$ and integral operators $q_{j}{ }^{*}$ in the sense adopted in [1], $t$ is time, and $q(t-\tau)$ is a kernel of the operator $q^{*}$; depending on the difference of the arguments.

If $\varphi$ is a rational function of the resolvent operators of the same class, then the expression ( 1,1 ) reduces to quadratures according to the well known rules [1]. In the general case the numerical realization of (1.1) is achieved by writing the function $\varphi$ in the form of a series in powers of the operators $q^{*}$, and then applying the formulas (1.1). A detailed interpretation of the functions of the operators and consequent computations are, as a rule, very time-consuming except in the case of small $t$, the latter ensuring the rapid convergence of the series. It is therefore expedient to construct effective computational algorithms for the expressions (1.1).

When Laplace transforms are applied to the initial relationships of the theory of hereditary elasticity we find, that in order to compute (1.1), it is necessary to invert the expression $\varphi(x ; Q(p)) F(p)$ where $p$ is the transformation parameter, $Q(p)=$
$L\{q(t)\} \quad$ and $F(p)=L\{f(t)\}$ are the Laplace transforms of the kernel $q(t)$ of the operator $q^{*}$ and of the function $f(t)$ respectively. Applying this to the operators $q^{*}$ with the hereditary kernels appearing as the fractional power exponents $\ni_{\alpha}$ ( $\beta$,
$t$ ) due to Rabotnov [1] and the function $P_{\alpha}(\lambda, t)$ due to Rzhanitsyn [2]

$$
\begin{equation*}
\ni_{\alpha}(\beta, t)=\sum_{k=0}^{\infty} \frac{\beta^{k} t^{k r+\alpha}}{\Gamma[r(k+1)]}, \quad \beta<0, \quad 0<r=1+\alpha<1 \tag{1.2}
\end{equation*}
$$

$$
\begin{equation*}
P_{\alpha}(\lambda, t)=\frac{1}{\Gamma(r)} t^{\alpha} e^{\lambda,}, \quad \lambda \leqslant 0, \quad 0<r=1+\alpha<1 \tag{1.3}
\end{equation*}
$$

the Laplace transforms of which are analytic functions with a characteristic singularity in the form of a branch point on the complex $p$-plane, we can compute the convolutions of the type (1.1) over a wide range of values of $t$.
2. Let $\varphi$ be a function of a single opeator $q^{*}$ and

$$
\begin{equation*}
Q(p)=Q_{1}(z), z=(p-\lambda)^{r}, \lambda \leqslant 0,0<r<1 \tag{2.1}
\end{equation*}
$$

According to the Mellin formula,

$$
\begin{align*}
& \varphi\left(q^{*}\right) f(t)=\frac{1}{2 \pi i} \int_{\sigma-i \infty}^{\sigma+i \infty} \varphi(Q(p)) F(p) e^{p t} d p  \tag{2.2}\\
& \varphi\left(q^{*}\right) \equiv \varphi\left(x ; q^{*}\right), \varphi(Q(p)) \equiv \varphi(x ; Q(p))
\end{align*}
$$

Here the integration is carried out along the straight line $\operatorname{Re} p=\sigma$ situated in the $p$-plane to the right of all singularities of the integrand function. We shall assume that in the finite part of the plane, with $|\arg p| \leqslant \pi$, this function can only have a single branch point $p=\lambda$ and, perhaps, a finite number of poles, none of them lying on the negative real semi-axis to the left of the branch point. To separate a single-valued branch of the function, we produce a cut along this semi-axis to the left of the point $p=\lambda$. Taking into account the residues of the poles, we replace the integration path in (2.2) by a contour consisting of the upper $\Gamma_{+}$and lower $\Gamma_{-}$ edge of the cut, the arcs $\Gamma_{1}$ and $\Gamma_{2}$ of the circle of infinite radius and of the circle $\Gamma_{\varepsilon}$ of vanishingly small radius, with centers at the point $p=\lambda$. Assuming that the conditions of the Jordan lemma hold and the integrals along the arcs $\Gamma_{1}$ and $\Gamma_{2}$ vanish, we obtain

$$
\begin{equation*}
\varphi\left(q^{*}\right) f(t)=\frac{1}{2 \pi i} \int_{H} \varphi(Q(p)) F(p) e^{p t} d p+\sum_{s} \operatorname{Res}\left[\varphi(Q(p)) F(p) e^{p t}\right] \tag{2.3}
\end{equation*}
$$

Here the contour $H=\Gamma_{-}+\Gamma_{\varepsilon}+\Gamma_{+}$is traversed in the anticlockwise direction, and the residues are calculated over all poles $p_{s},\left|\arg \left(p_{s}-\lambda\right)\right| \neq \pi$.

Substitution of $z=(p-\lambda)^{r}$ from (2.1) maps the contour $H$ onto the contour $H_{1}$ in the $\dot{z}$-plane, the latter contour consisting of the arcs $\arg z= \pm \pi r$ and the arc of the circle $|z|=\varepsilon^{r}=\varepsilon_{1}$. Let

$$
\varphi(Q(p))=\varphi_{1}(z)=z^{-v} \Phi(z)
$$

Here $v \geqslant 0$, and $\Phi(z)$ is a single-valued function analytic on the contour $H_{1}$ and in the region $|z| \leqslant \varepsilon_{1}$. In this case we have the following approximation [3] on $H_{1}$ :

$$
\begin{equation*}
\Phi(z)=\sum_{k=0}^{n} \Phi^{(k)}(0) \frac{z^{k}}{k!}+\rho_{n}(z), \quad \Phi^{(k)}(z)=\frac{d^{k} \Phi}{d z^{k}}, \quad n \geqslant 0 \tag{2.4}
\end{equation*}
$$

$$
\begin{equation*}
\left|\rho_{n}(z)\right| \leqslant M_{n+1}^{\prime} \frac{|z|^{n+1}}{(n+1)!}, \quad M_{n+1}^{\prime}=\max _{l}\left|\frac{d^{n+1}}{d z^{n+1}} \Phi(z)\right| \tag{2.5}
\end{equation*}
$$

where $l$ is a line connecting the point $z$ of the contour $H_{1}$ with the point $z=0$.
We assume that the following representation is possible for the function $F(p)$ analytic on $H$, but not necessarily regular at the point $p=\lambda$, for all $p \in H$ :

$$
\begin{array}{ll}
F(p)=\sum_{j=0}^{m} b_{j}(p-\lambda)^{i+\mu}+\gamma_{m}, & m \geqslant 0, \quad \mu \gtrless 0 \\
\left|\gamma_{m}\right| \leqslant A|p-\lambda|^{m+\mu+1}, & A=\mathrm{const} \geqslant 0 \tag{2.7}
\end{array}
$$

The conditions (2.6) and (2.7) are satisfied; in particular, by the transforms of the functions $t^{\delta}(\delta>-1), \sin \omega t, \cos \omega t, e^{h t}(h \geqslant \lambda)$.

By virtue of the equation [4]

$$
\begin{equation*}
\frac{1}{2 \pi i} \int_{H}(p-\lambda)^{\omega} e^{p t} d p=T_{\mathrm{I}}(t ; \omega) e^{\lambda t}, \quad T_{\mathrm{I}}(t ; \omega) \equiv \frac{t^{-\omega-1}}{\Gamma(-\omega)}, \quad \omega \gtrless 0 \tag{2,8}
\end{equation*}
$$

the expansions (2.4) and (2.6) transform (2.3) to the form

$$
\begin{align*}
& \varphi\left(q^{*}\right) f(t)=e^{\lambda t} \sum_{k=0}^{n} \sum_{j=0}^{m} b_{j} \frac{\Phi^{(k)}(0)}{k!} T_{\mathbf{I}}(t ; r k+j+\mu-r v)+  \tag{2.9}\\
& \quad \sum_{s} \operatorname{Res}_{p_{s}}\left[\varphi(Q(p)) F(p) e^{p t}\right]+r_{m n}(t) \\
& r_{m n}(t)=  \tag{2.10}\\
& \quad \frac{1}{2 \pi i}\left(\int_{\Gamma_{-}}+\int_{\Gamma_{\mathrm{e}}}+\int_{\Gamma_{+}}\right)\left\{e^{p t}\left[\gamma_{m} \varphi(Q(p))+\rho_{n} \sum_{j=0}^{m} b_{j}(p-\lambda)^{j+\mu-r v}\right] d p\right\}
\end{align*}
$$

If the numbers $n$ and $m$ in (2.4) and (2.6) satisfy the inequalities

$$
\begin{equation*}
r(n+1)+\mu>r v-1, \quad m+\mu+1>r v-1 \tag{2.11}
\end{equation*}
$$

then the integral along $\Gamma_{\varepsilon}$ in (2.10) vanishes as $\varepsilon \rightarrow 0$ by virtue of (2.5) and (2.7) and the boundedness of the function $\Phi(z)$ near the point $z=0$. The remaining integrals along the edges $\Gamma_{+}$and $\Gamma_{-}$of the cut on which $p-\lambda=x e^{ \pm i \pi}, x=$ $|p-\lambda|$, yield the formula

$$
\begin{equation*}
r_{m n}(t)=-\frac{1}{\pi} e^{\lambda i} \int_{0}^{\infty} e^{-x t} V(x) d x \tag{2,12}
\end{equation*}
$$

Here $V(x)$ denotes the imaginary part of the expression within the square brackets in (2.10) at the upper edge of the cut, i.e.

$$
\begin{equation*}
V(x)=\operatorname{Im}\left[\gamma_{m} \varphi(Q(p))+\rho_{n} \sum_{j=0}^{m} b_{j}(p-\lambda)^{j+\mu-\tau v}\right]_{p=\lambda+x e^{i \pi}} \tag{2.13}
\end{equation*}
$$

We note that if $\rho_{n}=0$ or $\gamma_{m}=0$, then the first or second condition of (2.11), becomes respectively superfluous.

Taking into account the relations (2.5) and (2.7) we obtain, from(2.12) and (2.13), the following estimate (where the maximum is taken over $\arg z= \pm \pi r$ ):

$$
\begin{align*}
& \left|r_{m n}(t)\right| \leqslant \frac{1}{\pi} e^{\lambda t}\left\{A M_{0} T_{2}(t ; m+\mu-r v+2)+\right.  \tag{2.14}\\
& \left.\quad \frac{M_{n+1}}{(n+1)!} \sum_{j=0}^{m}\left|b_{j}\right| T_{2}(t ; r n+r+j+\mu-r v+1)\right\} \\
& T_{2}(t ; \omega) \equiv \Gamma(\omega) t^{-\omega_{+}}, \quad M_{0}=\max |\Phi(z)|, \quad M_{n+1}=\max \left|\frac{d^{n+1} \Phi(z)}{d z^{n+1}}\right|
\end{align*}
$$

When the values of $t$ are sufficiently large, then (2.14) implies that $r_{m n}(t)$ can be neglected and the expression (2.9) used as an asymptotic expansion.Moreover, the formulas (2.9) and (2.12) make possible the calculation of the convolution $\varphi\left(q^{*}\right) f(t)$ also for other values of $t$, provided that the quantities $\rho_{n}$ and $\gamma_{m}$ in (2.13) are expressed in terms of the relations (2.4) and (2.6) and an approximate value of the quadrature(2.12) is found. The numbers $n$ and $m$ should be chosen as small as possible, with (2.11) taken into account. The integral in (2.12) is a Laplace transform, with parameter $t$, of the function of real variable $V(x)$. This proves the following theorem.

Theorem. Let
(1) a function $\varphi(u)$ analytic in the circle $|u| \leqslant u_{0}\left(u_{0}=\right.$ const $\left.>0\right)$ define a function of the hereditary operator $\varphi\left(q^{*}\right)$ (1.1) and let the Laplace transform $Q(p)$ of the operator $q^{*}$ on the plane of the transformation parameter have a branch point $p=\lambda(\lambda \leqslant 0)$ and $\varphi(Q(p))=(p-\lambda)^{-r \omega} \Phi\left[(p-\lambda)^{r}\right]$, $0<r<1, v \geqslant 0$ where $\Phi(z)$ is a single-valued function on the rays arg $z-$ $\pm \pi r$, at the point $z=0$ and in its neighborhood;
(2) the operator function $\varphi\left(q^{*}\right)$ act on the function $f(t)$ the Laplace transform $F(p)$ of which admits, on the negative semiaxis, at $p<\lambda$ and near the point $p=\lambda$, the representation $(2.6),(2.7)$ :
(3) the product $\varphi(Q(p)) F(p)$ tend to zero uniformly in $\arg p$ as $p \rightarrow \infty$ and $\operatorname{Re} p<\sigma=$ const ;
(4) amongst the singularities of the expression $\varphi(Q(p)) F(p)$ there exists, on a finite part of the $p$-plane with $|\arg p| \leqslant \pi$, only a single branch point $p=\lambda$ and possibly a finite number of poles none of which lie on the negative part of the real semi-axis to the left of the point $p=\lambda$.

Then the convolution $\varphi\left(q^{*}\right) f(t)$ will be given, for any finite $t>0$, by the formulas (2.9), (2.12) and (2.13), and at sufficiently large $t$ the relations (2.9) will assume the meaning of an asymptotic equation with (2.14) providing the estimate of its remainder term.

The formula (2.9) enables us, in particular, to establish the behavior of the convolution $\varphi\left(q^{*}\right) f(t)$ as $t \rightarrow \infty$. A finite limit value exists if $\operatorname{Re} p_{s}<0$ and
either $\lambda<0$ or $\lambda=0$ for all poles $p_{s}$. at which the residues are computed without however positive powers of $t$ appearing in the expansion (2.9). When $\lambda=$ 0 and $f(t)=$ const , the index $\mu=-1$ and, provided that the conditions Reps $<0, \Phi(0) \neq 0$ hold, the convolution is bounded if $v=0$.

From the known theorem of operational calculus we have

$$
\begin{align*}
& \lim _{t \rightarrow \infty} q^{*} \cdot 1=\lim _{p \rightarrow 0} p\left[Q(p) \frac{1}{p}\right]=Q(0)  \tag{2.15}\\
& \lim _{t \rightarrow \infty} \varphi\left(q^{*}\right) f(t)=\lim _{p \rightarrow 0} p \varphi(Q(p)) F(p)=\varphi(Q(0)) \lim _{p \rightarrow 0} p F(p)= \\
& \quad \varphi\left(\lim _{t \rightarrow \infty} q^{*} \cdot 1\right) f(\infty)
\end{align*}
$$

provided that the limits in question exist. The assertion (2.15) represents a limiting theorem for the hereditary operators, established in [1].
3. Let us now consider the operator $q^{*}$ with the kernel (1.2) $x \ni_{\alpha}(\beta, t), x>0$. Since $L\left\{\vartheta_{\alpha}(\beta, t)\right\}=\left(p^{r}-\beta\right)^{-1}$, we must put $\lambda=0$ and $z=p$ and assume, in the case when $\varphi(Q(p))$ has no singularity at the point $p=0$, that $v=0$. We also have

$$
\begin{aligned}
& \Phi(z)=\varphi(x /(z-\beta)), \Phi^{(k)}(0)=(-1)^{k} \partial^{k} \varphi(-x / \beta) / \partial \beta^{k} \\
& \varphi\left(x \vartheta_{\alpha}^{*}(\beta)\right) f(t)=\sum_{k=0}^{n} \sum_{j=0}^{m} b_{j} \frac{(-1)^{k}}{k!} \frac{\partial^{k} \varphi(-x / \beta)}{\partial \beta^{k}} T_{\mathrm{I}}(t ; r k+j+\mu)+ \\
& \quad \sum_{s} \operatorname{Res}_{p_{s}}\left[\varphi\left(\frac{\chi}{p^{r}-\beta}\right) F(p) e^{p t}\right]+r_{m n}(t)
\end{aligned}
$$

Let the function $\varphi(u)$ be analytic in the circle $|u| \leqslant u_{0}$ and

$$
\max _{\arg z= \pm \pi r}\left|\frac{x}{z-\beta}\right|=\frac{x}{g}<u_{0}
$$

Then, when $r \leqslant 1 / 2 \quad g=|\beta|$ and $r>1 / 2 \quad g=|\beta| \sin \pi r$. In this case the quantities $M_{0}$ and $M_{n+1}$ from (2.14) can be written in the form of a series. In particular, for $M_{n+1}$ we have

$$
\begin{aligned}
& M_{n+1}=\max _{\arg z= \pm \pi r}\left|\sum_{j=0}^{\infty}(-1)^{n+1} S_{j n}(z-\beta)^{-j-n-1}\right| \\
& S_{j n} \equiv(j+n)!x^{j} \varphi^{(j)}(0) /[(j-1)!j!]
\end{aligned}
$$

If all $\varphi^{(j)}(0), j=0,1,2, \ldots$ have the same sign, then

$$
\begin{equation*}
M_{n+1} \leqslant\left|\sum_{j=0}^{\infty} S_{j n} g^{-j-n-1}\right|=\left|\frac{\partial^{n+1} \varphi(x / g)}{\partial g^{n+1}}\right|, \quad M_{0} \leqslant\left|\varphi\left(\frac{x}{g}\right)\right| \tag{3.2}
\end{equation*}
$$

When the sign alternates, i. e., $\operatorname{sign} \varphi^{(j)}(0)=-\operatorname{sign} \varphi^{(j+1)}(0), \quad j=0,1$, 2 , . . ., we have

$$
\begin{equation*}
M_{n+1} \leqslant\left|\sum_{j=0}^{\infty}(-1)^{j} S_{j n} g^{-j-n-1}\right|=\left|\frac{\partial^{n+1} \varphi(-x / g)}{\partial g^{n+1}}\right|, \quad M_{0} \leqslant\left|\varphi\left(-\frac{\chi}{g}\right)\right| \tag{3.3}
\end{equation*}
$$

The above inequalities simplify the estimation of $r_{m n}(t)$ when formula (2.14) is used.

When $f(t)=$ const and $\lambda=0$ in (2.6), we have $b_{0} \neq 0, b_{j}=0(j \geqslant$ 1), $\mu=-1, \gamma_{m}=0$, and the formula (3.1) coincides, in the absence of the residues, with the asymptotic expansion obtained in [5]. If $f(t)=t^{\circ}(\delta>-1)$, $\varphi\left(q^{*}\right)=q^{* \omega}, \omega=1,2, \ldots, q(t)=x \ni_{\alpha}(\beta, t)$, then $b_{0}=\Gamma(1+\delta)$, $b_{j}=0(j \geqslant 1), \mu=-\delta-1, \gamma_{m}=0$. In this case the formula (3.1) and the estimates (2.14) and (3.2) yield the relations obtained in [6].

Example 1. Let $\varphi\left(x \partial_{\alpha}^{*}(\beta)\right)=\left(1+x \partial_{\alpha}{ }^{*}(\beta)\right)^{1 / 2}, x>0, \beta<0, f(t)=1$. Then $F(p)=p^{-1}$ and in (2.6) we have $\lambda=0, b_{0}=1, m=0, \mu=-1, \gamma_{m}=0$. When $n=1$, we have, in accordance with (3.1),

$$
\left(1+x \ni_{\alpha} *(\beta)\right)^{1 / 2} \cdot 1=\left(1-\frac{x}{\beta}\right)^{1 / 2}-\frac{x}{2 \beta^{2}\left(1-x \beta^{-1}\right)^{1 / 2}} \frac{t^{-r}}{\Gamma(1-r)}+r_{1}(t)
$$

From (2.14) we obtain

$$
\left|r_{1}(t)\right| \leqslant M_{2} \Gamma(2 r) t^{-2 r} /(2 \pi), \quad M_{2} \leqslant x g^{-5 / 2} g_{1} g^{-1 / 2}\left(x /\left(4 g_{1}\right)+1\right)
$$

Here, when $r \leqslant 1 / 2 g=|\beta|, g_{1}=\left|\beta_{1}\right|$, and when $r>1 / 2 g=|\beta| \sin \pi r, g_{1}=\left|\beta_{1}\right|$
$\sin \pi r, \quad \beta_{1}=\beta-\chi$. Thus, if $\alpha=-0.7(r=0.3), \beta=-1, x=0.5$, then $\left(1+x \ni_{\alpha}{ }^{*}(\beta)\right)^{1 / 2} \cdot 1=1.225-0.157 t^{-0.3}+r_{1}(t)$ and $\left|r_{1}(t)\right| \leqslant 0.105 t^{-0.6}$. The above results can be extended to the case of two and more operator arguments of the function $\varphi$. In particular, when $\varphi=\varphi\left(\chi_{1} \ni_{\alpha_{1}}{ }^{*}\left(\beta_{1}\right), \chi_{2} \ni_{\alpha_{2}}{ }^{*}\left(\beta_{2}\right)\right)$ and under the condition that the operational analog of this function is bounded in the neighborhood of the point $p=0$ and $1+\alpha_{1}=c_{1} r, 1+\alpha_{2}=c_{2} r$ where $c_{1}$ and $c_{2}$ are positive integers and $0<r<1$, we have the following relations for the function $\Phi(z)$ :

$$
\Phi(z)=\varphi\left(x_{1} /\left(z^{c_{1}}-\beta_{1}\right), \quad x_{2} /\left(z^{c_{t}}-\beta_{2}\right)\right), \quad z=p^{r}
$$

Using similar arguments, we obtain

$$
\sum_{s} \operatorname{Res}\left[\varphi\left(\frac{x_{1}}{p^{c_{1} r}-\beta_{1}}, \frac{x_{2}}{p^{c_{2} r}-\beta_{2}}\right) F(p) e^{p t}\right]+r_{m n}(t)
$$

When $c_{1}=c_{2}=1$, the following relation holds:

$$
\Phi^{(k)}(0)=(-1)^{k}\left(\partial / \partial \beta_{1}+\partial / \partial \beta_{2}\right)^{k} \varphi\left(-x_{1} / \beta_{1},-x_{2} / \beta_{2}\right)
$$

The estimate of $r_{m n}(t)$ retains the form of (2.14) when $v=0, \lambda=0$.
Example 2. The operator expression $\left(1+x_{1} \vartheta_{\alpha}^{*}\left(\beta_{1}\right)+x_{2} \vartheta_{\alpha}^{*}\left(\beta_{2}\right)\right)^{1 / 2} \cdot 1, x_{1}$, $x_{2}>0, \beta_{1}, \beta_{2}<0$, has, according to (3.4), the following representation:

$$
\left(1+x_{1} \ni_{\alpha} *\left(\beta_{1}\right)+x_{2} \ni_{\alpha} *\left(\beta_{2}\right)\right)^{1 / 2} \cdot 1=\left(1-\frac{x_{1}}{\beta_{1}}-\frac{x_{2}}{\beta_{2}}\right)^{1 / 2}-
$$

$$
\frac{x_{1} \beta_{1}^{-2}+x_{2} \beta_{2}^{-2}}{2\left(1-x_{1} \beta_{1}^{-1}-x_{2} \beta_{2}^{-1}\right)^{1 / 2}} \frac{t^{-r}}{\Gamma(1-r)}+r_{1}(t)
$$

From (2.14) we obtain

$$
\begin{gathered}
\left|r_{1}(t)\right| \leqslant \frac{M_{2} \Gamma(2 r)}{2 \pi} t^{-2 r}, \quad M_{2} \leqslant\left(1+\frac{\left|\beta_{1}-\beta_{2}\right|}{g}\right)^{1 / 2} \times \\
\left(\frac{x^{2}}{g^{4}}+\frac{x^{2}}{2 g^{5}}\left|\beta_{1}-\beta_{2}\right|+\frac{2 x}{g^{3}}\right) \\
r=1+\alpha, \quad x=\max \left\{x_{1}, x_{2}\right\}, \quad g=\min \left\{g_{j}\right\}, \quad j=1,2,3,4 ; \quad r \leqslant 1 / 2 \quad g_{j}=\left|\beta_{j}\right| \\
r>1 / 2 \quad g_{j}=\left|\beta_{j}\right| \sin \pi r ; \quad \beta_{3,4}=-1 / 2\left[e_{1}+e_{2} \pm\left(\left(e_{1}-e_{2}\right)^{2}+4 x_{1} x_{2}\right)^{2 / 2}\right]
\end{gathered}
$$

$$
\left(e_{1}=x_{1}-\beta_{1}, e_{2}=x_{2}-\beta_{2}\right)
$$

Asymptotic expansion of the function of the operator $q^{*}$ with kernel (1.3) is obtained in the analogous manner.

Example 3. Let $\varphi\left(q^{*}\right) f(t)=q^{* \omega} \cdot 1, \omega=1,2, \ldots, q(t)=x P_{\alpha}(\lambda, t), x>$ $0, \lambda<0$. Then

$$
\begin{aligned}
& \varphi(Q(p))=x^{\omega} z^{-\omega}, \quad z=(p-\lambda)^{r}, \quad \Phi(z)=\Phi(0)=x^{\omega}, \quad v=\omega \\
& F(p)=\frac{1}{p}=\sum_{j=0}^{m}(-1)^{j} \lambda^{-j-1}(p-\lambda)^{j}+(-1)^{m+1} p^{-1} \lambda^{-m-1}(p-\lambda)^{m+1}
\end{aligned}
$$

Thus we have, in accordance with the expressions (2.4), (2.6), (2.7) and (2.14), $n=$ $0, p_{n}=0, \mu=0, A=|\lambda|^{-m-2}, M_{0}=\chi^{\omega}, M_{n_{巾 1}}=0$, and from (2.9), (2.14) follows

$$
\begin{align*}
& x^{\omega} P_{\alpha}^{*^{\omega}}(\lambda) \cdot 1=x^{\omega} e^{\lambda t} \sum_{j=0}^{m}(-1)^{j} \lambda^{-j-1} T_{1}(t ; j-r \omega)+  \tag{3.5}\\
& \quad \frac{x^{\omega}}{(-\lambda)^{r \omega}}+r_{m n}(t) \\
& \left|r_{m n}(t)\right| \leqslant x^{\omega}|\lambda|^{-m-2} e^{\lambda t} T_{2}(t ; m-r \omega+2) / \pi, \quad m-r \omega+2>0
\end{align*}
$$

If $r \omega$ is an integer, then the branch point of the function

$$
x^{\omega} /\left[p(p-\lambda)^{r \omega}\right]=L\left\{\chi^{\omega} P_{\alpha}^{* \omega}(\lambda) \cdot 1\right\}
$$

at $p=\lambda$ is replaced by a pole, and this simplifies the expression (3.5) since in this case $r_{m n}(t)=0$ and we have $T_{1}(t ; j-r \omega)=0$ when $j \geqslant r \omega$.
4. Using (2.12), we construct an analytic expression for approximating the quantity $r_{m n}(t)$. Let us approximate the function $V(x)$ using combinations of the exponential expressions and power functions. Since the multiplier $e^{-x t}$ decays rapidly, it is sufficient to attain a good approximation to $V(x)$ at the initial part of the interval of integration, ensuring at large $x$ only that the manner of behavior is similar.

Let

$$
V(x) \approx \sum_{j} a_{j} x^{\delta_{j} e^{\lambda_{j} x}}, \quad \delta_{j}>-1, \quad \lambda_{j} \leqslant 0
$$

Then

$$
\int_{0}^{\infty} e^{-x t} V(x) d x=\sum_{j} u_{j} \frac{\Gamma\left(\delta_{j}+1\right)}{(t-\lambda)^{\delta_{j}+1}}+\int_{0}^{\infty} e^{-x t}\left[V(x)-\sum_{j} a_{j} x^{\delta_{j} e^{\lambda_{j} x}}\right] d x
$$

The last integral can be used to assess the error of approximation.
To illustrate this, let us consider the convolution $\ni_{\alpha}{ }^{*}(\beta) \cdot 1$ the Laplace transform of which $\left(p^{r}-\beta\right)^{-1} p^{-1}$ for $0<r<1, \beta<0,|\arg p| \leqslant \pi$ has a unique singularity, namely a branch point at $p=0$. Since $\lambda=0$ and $F(p)=p^{-1}$, therefore we have in (2.6) $\quad b_{0}=1, m=0, \mu=-1$ and $\gamma_{m}=0$. Here $\Phi(z)=(z-$ $\beta)^{-1}$ and $v=0$, and to fulfil the conditions (2.11) it is sufficient to put $n=0$, which yields $\rho_{n}=\beta^{-1} p^{r}\left(p^{r}-\beta\right)^{-1}$. In accordance with (2.9), (2.12) and (2.13), we obtain [7]

$$
\begin{equation*}
\vartheta_{\alpha} *(\beta) \cdot 1=-\frac{1}{\beta}+\frac{\sin \pi r}{\pi \beta} \int_{0}^{\infty} e^{-x|\beta|^{1 / r} t} \frac{x^{r-1} d x}{x^{2 r}+2 x^{r} \cos \pi r+1} \tag{4.1}
\end{equation*}
$$

Let $r=0.3(\alpha=-0.7)$. We interpolate the decaying function $x^{r-1}\left(x^{2 r}+2 x^{r}\right.$ $\cos \pi r+1)^{-1}$ on the interval $[0,4]$ using the expression

$$
x^{r-1} e^{-0.3 x}\left(1-1.18 x^{r}+0.587 x^{2 r}+0.02 x^{3 r}\right)
$$

From (4. 1) follows

$$
\begin{gather*}
\exists_{-0.7}^{*}(\beta) \cdot 1 \approx-\beta^{-1}+(\pi \beta)^{-1} \sin \pi r \sum_{j=1}^{4} A_{j} \Gamma^{\Gamma}(i r) y^{j} \quad A_{1}=1  \tag{4.2}\\
A_{2}=-1.18, \quad A_{3}=0.587, \quad A_{4}=0.02, \quad y=(\tau+0.3)^{-r}, \quad \tau=|\beta|^{1 / r} t
\end{gather*}
$$

Comparison of the values computed with help of the formula (4.2) with those given in tables [8] shows that at $\tau=0.1$ the error is about $5 \%$, at $\tau=0.2$ it does not exceed $1.5 \%$ and at $\tau \geqslant 0.3$ the values practically coincide. It must however be remembered that the asymptotic expansion of $\ni_{\alpha}{ }^{*}(\beta) \cdot 1$ [8] hold only when $\tau \geqslant 1$.

## REFERENCES

1. Rabotnov, Iu. N. Elements of the Hereditary Mechanics of Solids. Moscow, "Nauka", 1977.
2. Rzhanitsyn, A. R. Theory of Creep. Moscow, Stroiizdat, 1968.
3. Goncharov, V. L. Theory of Interpolation and Approximation of Functions. Moscow, Gostekhizdat, 1954.
4. Bateman, H. and Erdelyi, A. Higher Transeendental Functions. Vol. 1, McGraw - Hill, New York, 1953.
5. Dolinina, N. N. On the functions of special operators of the theory of elastic media with memory. Dokl. Akad. Nauk SSSR, Vol. 170, No. 1, 1966.
6. Sinaiskin, E. S. On the asymptotic representation of the operator used in describing the behavior of the elastic media with memory. Akad. Nauk SSSR, Mekhanika, No. 1, 1965.
7. Listovnichii, V. F. and Shermergor, T. D. Creep of viscoelastic media with kernel of the type of a degenerate hypergeometric function. Izv. Akad. Nauk SSSR, MTT, No. 1, 1969.
8. Rabotnov, Iu. N., Papernik, L. Kh. and Zvonov, E. N. Tables of Fractionally exponential Function of Negative Parameters, and its Integral. Moscow, "Nauka", 1969.
